

DEFORMATION RIGIDITY OF SOME SIMPLE AFFINE VOAS

ANDREW R. LINSHAW AND FEI QI

ABSTRACT. In this paper, we prove that simple affine vertex operator algebras with positive integer levels admit only trivial first-order deformations. Therefore, the deformation rigidity conjecture of strongly rational vertex operator algebras holds for these cases. We also show that the same holds for simple affine vertex operator algebra of $\mathfrak{sl}_2$ at the non-integral admissible level $-4/3$. Therefore, neither C_2 -cofiniteness nor rationality is a necessary condition for deformation rigidity of VOAs. We conjecture that the same should hold for every simple affine VOA that does not coincide with the corresponding universal affine VOA.

1. INTRODUCTION

The moduli space of conformal field theories is very important in both physics and mathematics and connects to a number of branches of mathematics. But we know very little about this moduli space, not even the correct topology. It has been a common belief among physicists since the 1990s that among all two-dimensional conformal field theories, the strongly rational ones are isolated. The belief is formulated into a mathematical conjecture by Yi-Zhi Huang in [H4], Problems 4.1 and 4.2. It is therefore important to study the deformations of vertex algebras. One of the conjectures in this study is what we call the deformation rigidity conjecture: the deformation of a rational vertex operator algebra that preserves the grading and the vacuum can only be trivial.

Classifying the first-order deformations has been a long-standing challenge, mainly due to technical convergence conditions. The first significant result did not appear until 2024, when the second author, in collaboration with Vladimir Kovalchuk, overcame the difficulty in convergence, and gave a programmable algorithm for classifying the first-order deformations for every freely generated vertex algebra (see [KQ]).

One of the key theorems in [KQ] states that for the vertex operator Y_1 associated with the first-order deformation, if for any generators a, b of the vertex algebra, $Y_1(a, x)b$ contains no negative powers, then Y_1 gives a trivial deformation. This theorem allows us to focus only on the singular part of $Y_1(a, x)b$ with only finitely many structural constants. [KQ] proved the key theorem by constructing the linear map corresponding to the trivial deformation based on the PBW-type basis on a freely generated vertex algebra. The proof fails if there are relations among the basis vectors.

In the case of simple affine vertex operator algebras $L_k(\mathfrak{g})$ with the level k being a positive integer level, we show that the same result holds by using a theorem of Bong Lian [L] to circumvent the direct construction in [KQ]. Then, by combining some of the computations in [KQ], we

finish the proof of the deformation rigidity conjecture for these particular examples of rational vertex operator algebras.

We also prove the same result for $L_{-4/3}(\mathfrak{sl}_2)$, a particular case of a simple affine VOA at a non-integral admissible level. Note that $L_{-4/3}(\mathfrak{sl}_2)$ is neither C_2 -cofinite nor rational, so neither of these conditions is necessary for deformation rigidity. We conjecture that the deformation rigidity should hold for all simple affine VOAs $L_k(\mathfrak{g})$ that do not coincide with the corresponding universal affine VOA $V^k(\mathfrak{g})$. If the conjecture holds, then together with the results in [KQ], we have

$$\dim H_{1/2}^2(L_k(\mathfrak{g}), L_k(\mathfrak{g})) = \begin{cases} 1 & \text{if } L_k(\mathfrak{g}) = V^k(\mathfrak{g}) \\ 0 & \text{if } L_k(\mathfrak{g}) \neq V^k(\mathfrak{g}) \end{cases}.$$

The proof, however, might require a completely different method.

The paper is organized as follows. Section 2 reviews the notion of first-order deformations and their equivalence classes. Section 3 shows that a regular Y_1 is trivial. In Section 4, we give a proof that Y_1 is regular for the simple affine VOA $L_k(\mathfrak{g})$ with $k \in \mathbb{Z}_+$. In Section 5, we give a proof that Y_1 is regular for the simple affine VOA $L_{-4/3}(\mathfrak{sl}_2)$. Both Section 4 and 5 substantially depend on the formulas and computations in Section 4.1 and 6.2 of [KQ].

Acknowledgement: The second author would like to thank Yi-Zhi Huang for his long-term support. The first author is supported by NSF Standard Grant DMS-2401382 and Simons Foundation Grant MPS-TSM-00007694. Both authors would like to thank Yi-Zhi Huang, Shashank Kanade, and Paul Johnson for their comments on the early versions of the draft.

2. PRELIMINARIES

2.1. First-order deformations. Let $(V, Y, \mathbf{1})$ be a grading-restricted vertex algebra in the sense of [H1]. This means a $\mathbb{Z}$ -graded vertex algebra with usual grading restrictions, and with only the grading operator $\mathbf{d} = L(0)$ and $D = L(-1)$. For $u, v \in V$, we keep Borcherds' notation u_n for the modes of $Y(u, x)$, i.e.,

$$Y(u, x)v = \sum_{n \in \mathbb{Z}} u_n v x^{-n-1}$$

A first-order deformation of V is defined by a vertex operator

$$\begin{aligned} Y_1 : V \otimes V &\rightarrow V((x)) \\ u \otimes v &\mapsto Y_1(u, x)v = \sum_{n \in \mathbb{Z}} u_n^{def} v x^{-n-1}. \end{aligned}$$

such that the vector space $V^t = \mathbb{C}[t]/(t^2) \otimes_{\mathbb{C}} V$, with the vertex operator

$$Y^t(u, x)v = Y(u, x)v + tY_1(u, x)v,$$

and the vacuum element $\mathbf{1}$, forms a grading-restricted vertex algebra over the base ring $\mathbb{C}[t]/(t^2)$ (see [H2] and [KQ] for more details). Equivalently, Y_1 should satisfy the following conditions:

(1) For every $u, v \in V$, the mode u_n^{def} satisfies

$$\text{wt } u_n^{def} v = \text{wt } u - n - 1 + \text{wt } v.$$

(2) For every $v \in V$,

$$Y_1(\mathbf{1}, x)v = 0, Y_1(u, x)\mathbf{1} = 0. \quad (1)$$

In particular, the D -operator

$$D(u + tv) = \lim_{x \rightarrow 0} \frac{d}{dx} Y^t(u + tv, x)\mathbf{1} = \lim_{x \rightarrow 0} \frac{d}{dx} Y(u + tv, x)\mathbf{1} = u_{-2}\mathbf{1} + tv_{-2}\mathbf{1}$$

on V^t is defined purely with modes of Y , regardless of Y_1 .

(3) For every $u, v \in V$, the D -derivative-commutator formula holds,

$$[D, Y_1(u, x)]v = Y_1(Du, x)v = \frac{d}{dx} Y_1(u, x)v.$$

(4) For $u, v \in V$, the following Jacobi-like identity holds,

$$\begin{aligned} & x_0^{-1} \delta \left(\frac{x_1 - x_2}{x_0} \right) (Y_1(u, x_1)Y(v, x_2) + Y(u, x_1)Y_1(v, x_2)) \\ & - x_0^{-1} \delta \left(\frac{-x_2 + x_1}{x_0} \right) (Y_1(v, x_2)Y(u, x_1) + Y(v, x_2)Y_1(u, x_1)) \\ & = x_2^{-1} \delta \left(\frac{x_1 - x_0}{x_2} \right) (Y_1(Y(u, x_0)v, x_2) + Y(Y_1(u, x_0)v, x_2)). \end{aligned}$$

Remark 2.1. For $u, v \in V$ and $m, n \in \mathbb{Z}$, by taking an appropriate residue of the Jacobi-like identity, we have the following commutator condition.

$$[u_m^{def}, v_n] + [u_m, v_n^{def}] = \sum_{\alpha \geq 0} \binom{m}{\alpha} \left((u_\alpha^{def} v)_{m+n-\alpha} + (u_\alpha v)_{m+n-\alpha}^{def} \right). \quad (2)$$

The commutator relation (2) should hold for every $u, v \in V$ and $m, n \in \mathbb{Z}$. In case V is a freely generated vertex algebra, [KQ] proves, via a lengthy process, that it suffices for (2) to hold for every pair of generators u, v and every $m, n \in \mathbb{N}$. The result does not naively generalize to non-freely generated vertex algebras. But we would not need it if no nontrivial first-order deformations exist.

2.2. Equivalent first-order deformations. Two first-order deformations $(V^t, Y^{t,(1)})$ and $(V^t, Y^{t,(2)})$ are equivalent, if there exists a $\mathbb{C}[t]/(t^2)$ -linear vertex algebraic isomorphism $f^t : V \oplus tV = V^t \rightarrow V^t = V \oplus tV$ whose restriction on V is of the form

$$f^t|_V = 1_V + tf_1,$$

where $f_1 : V \rightarrow V$ is a $\mathbb{C}$ -linear grading-preserving map. In other words, for $u, v \in V$,

$$f^t(u + tv) = u + t(v + f_1(u)).$$

In the papers [H1] and [H2], Huang showed that the equivalence classes of first-order deformations are described by the second cohomology $H_{1/2}^2(V, V)$ of vertex algebras. Here we shall not give a detailed review of the definitions of 2-cocycles and 2-coboundaries, but only mention one

consequence. Fix an $Y_1 : V \otimes V \rightarrow V((x))$. If there exists a grading preserving map $\phi : V \rightarrow V$ such that $\phi(\mathbf{1}) = 0$ (see [H3]), and for every $u, v \in V$,

$$Y_1(u, x)v = Y(\phi(u), x)v - \phi(Y(u, x)v) + Y(u, x)\phi(v),$$

then the first-order deformation given by Y_1 is equivalent to that given by 0. In this case, we say Y_1 is a trivial first-order deformation.

2.3. Simple affine VOA. In this paper, we focus on the particular example where V is the simple affine VOA (see [FZ], [LL] for more details). Let $\mathfrak{g}$ be a finite-dimensional simple Lie algebra with root system Φ . Let θ be the highest root of $\mathfrak{g}$. For every $\alpha \in \Phi$, let $e_\alpha, f_\alpha, h_\alpha$ be the basis vectors of the $\mathfrak{sl}_2$ -subalgebra of $\mathfrak{g}$ associated with α . When $\alpha = \theta$, we abbreviate $e_\theta, f_\theta, h_\theta$ as e, f, h . For a fixed real number k , let $V^k(\mathfrak{g})$ be the universal affine VOA associated with $\mathfrak{g}$ with level k . Let $L_k(\mathfrak{g})$ be the simple quotient of $V^k(\mathfrak{g})$. In case $k + h^\vee \in \mathbb{Q}_{\geq 0} \setminus \{1/m : m \in \mathbb{Z}_+\}$, $L_k(\mathfrak{g})$ is a nontrivial quotient of $V^k(\mathfrak{g})$ (see [GK] for details). In case k is admissible, $L_k(\mathfrak{g})$ is the quotient of $V^k(\mathfrak{g})$ by an ideal generated by singular vectors. (see [KW]). In case k is a positive integer, $L_k(\mathfrak{g})$ is the quotient of $V^k(\mathfrak{g})$ by the ideal generated by $e(-1)^{k+1}\mathbf{1}$ (see [LL]).

Hereafter, we set $V = L_k(\mathfrak{g})$, $V^t = \mathbb{C}[t]/(t^2) \otimes_{\mathbb{C}} V = V \oplus tV$ the first-order deformation of V with the vertex operator $Y^t = Y + tY_1$. For $a \in \mathfrak{g}$, we use $a(m)$ and $a^{def}(m)$ to denote the modes in the vertex operators $Y(a, x)$ and $Y_1(a, x)$, respectively.

3. TRIVIALITY OF A REGULAR Y_1

For a freely generated vertex algebra, [KQ] shows that if the Y_1 -operator on every pair of generators is regular, then Y_1 defines a trivial first-order deformation. Here, using a result in [L], we generalize it to the simple affine VOAs that are not freely generated.

Theorem 3.1. Let V be a simple affine VOA. If for every $a, b \in \mathfrak{g}$, $Y_1(a(-1)\mathbf{1}, x)b(-1)\mathbf{1}$ is regular, i.e., it contains no negative powers, then Y_1 defines a trivial first-order deformation.

Proof. Consider the $\mathbb{C}$ -subalgebra W of V^t generated by the vacuum $\mathbf{1}$ and the weight-1 subspace $V_{(1)} \subset V_{(1)} \oplus tV_{(1)} \subset V^t$. Since Y_1 is regular, for $a, b \in V_{(1)}$, the singular part of $Y^t(a, x)b$ coincides with that of $Y(a, x)b$. From Theorem 4.11 of [L], W is isomorphic to a quotient of affine VOA. Clearly, $W_{(1)} = V_{(1)}$ is isomorphic to $\mathfrak{g}$. So W is isomorphic to a quotient of $V^k(\mathfrak{g})$. In particular, $W_{(0)} = \mathbb{C}\mathbf{1}$.

Consider the projection $\pi : V^t \rightarrow V$ modulo t , i.e., $\pi(v + tv^{(1)}) = v$. Clearly, π is a vertex algebra homomorphism. Let $\pi|_W$ be its restriction on W . Since $W_{(1)} = V_{(1)}$, $\pi|_W$ is surjective. Let $K = \ker \pi|_W$. Clearly, K is an ideal of W contained in $tV \cap W$. Let

$$I = \{v \in V : tv \in K\}.$$

We show that I is an ideal of V . Fix any $u \in V$ and $v \in I$. Let $w = u + tw^{(1)} \in W$ so that $\pi(w) = u$. Since tv is an element of the ideal $K \subset W$, for every $n \in \mathbb{Z}$, $(w_n + tw_n^{def}) \cdot tv \in K$,

i.e.,

$$(w_n + tw_n^{def}) \cdot tv = (u_n + tw_n^{(1)} + tw_n^{def}) \cdot tv = tu_nv$$

is an element in K . Thus, $u_nv \in I$.

Since V is simple, either $I = 0$ or $I = V$. If $I = V$, then in particular, $\mathbf{1} \in I$, i.e., $t\mathbf{1} \in W$, contradicting $W_{(0)} = \mathbb{C}\mathbf{1}$. Therefore, $I = 0$. Consequently, $K = 0$. So $\pi|_W$ is an isomorphism. We thus proved that W is isomorphic to $V = L_k(\mathfrak{g})$.

Let $f = (\pi|_W)^{-1} : V \rightarrow W$ be the isomorphism. From $\pi \circ f = 1_V$, we know that for every $v \in V$

$$f(v) = v - t\phi(v)$$

where $\phi : V \rightarrow V$ is a homogeneous linear map. Fix any $u, v \in V$. We read the coefficient of t from

$$f(Y(u, x)v) = Y^t(f(u), x)f(v),$$

to obtain

$$-\phi(Y(u, x)v) = Y_1(u, x)v - Y(\phi(u), x)v - Y(u, x)\phi(v).$$

Therefore, $Y_1(u, x)v$ coincides with the trivial deformation given by ϕ , thus trivial. $\square$

4. THE POSITIVE INTEGER LEVEL CASE

In this section we show that Y_1 is regular. The computation in Section 6.2 of [KQ] generalizes to the current situation, showing that the first-order pole of Y_1 is represented by a trivial first-order deformation, and for every $a, b \in \mathfrak{g}$,

$$Y_1(a(-1)\mathbf{1}, x)b(-1)\mathbf{1} = c \cdot \langle a, b \rangle \mathbf{1}x^{-2} + \sum_{m \geq 1} a^{def}(-m)b(-1)\mathbf{1}x^{m-1} \quad (3)$$

for some constant $c \in \mathbb{C}$ (which is denoted by $B_{\mathfrak{g}}$ in [KQ]). Here $\langle \cdot, \cdot \rangle$ is the invariant bilinear form on $\mathfrak{g}$. So, it suffices to show that $c = 0$.

The key idea is to consider the coefficient of $e(-1)^k\mathbf{1}$ in the expression of $f^{def}(1)e(-1)^{k+1}\mathbf{1}$. This expression is zero because $e(-1)^{k+1}\mathbf{1} = 0$ in V . On the other hand, we may use the commutator condition (2) to calculate it. This shall establish an equation of c .

4.1. The mode $f^{def}(1)$. We now consider the commutator condition (2) with $u = f(-1)\mathbf{1}, v = e(-1)\mathbf{1}$, $m = 1$, and $n = -1$. Then the commutator condition reads

$$\begin{aligned} & [f^{def}(1), e(-1)] + [f(1), e^{def}(-1)] \\ &= \left(f^{def}(0)e(-1)\mathbf{1} \right)_{1-1-0} + (f(0)e(-1)\mathbf{1})_{1-1-0}^{def} + \left(f^{def}(1)e(-1)\mathbf{1} \right)_{1-1-1} + (f(1)e(-1)\mathbf{1})_{1-1-1}^{def} \\ &= 0 - h^{def}(0) + (c \cdot \mathbf{1})_{-1} + (k \cdot \mathbf{1})_{-1}^{def} = -h^{def}(0) + c. \end{aligned}$$

Unraveling the left-hand-side, we have

$$f^{def}(1)e(-1) - e(-1)f^{def}(1) + f(1)e^{def}(-1) - e^{def}(-1)f(1) = -h^{def}(0) + c. \quad (4)$$

Let $i \in \mathbb{Z}_+, 1 \leq i \leq k+1$. Apply (4) to $e(-1)^{i-1}\mathbf{1}$, we see that

$$\begin{aligned} f^{def}(1)e(-1)^i\mathbf{1} &= e(-1)f^{def}(1)e(-1)^{i-1}\mathbf{1} - f(1)e^{def}(-1)e(-1)^{i-1}\mathbf{1} + e^{def}(-1)f(1)e(-1)^{i-1}\mathbf{1} \\ &\quad - h^{def}(0)e(-1)^{i-1}\mathbf{1} + c \cdot e(-1)^{i-1}\mathbf{1}. \end{aligned} \quad (5)$$

In particular, if $i = k+1$, then left-hand-side is zero.

4.2. The term with $h^{def}(0)$. We start by analyzing $h^{def}(0)e(-1)^i\mathbf{1}$ and show that it is zero for every $1 \leq i \leq k$.

We argue by induction. From (3), we see that

$$h^{def}(0)e(-1)\mathbf{1} = 0.$$

Thus, the conclusion holds for $i = 1$.

Assume the conclusion holds for smaller i . In the commutator relation (2), we set $u = h(-1)\mathbf{1}, v = e(-1)\mathbf{1}, m = 0$, and $n = -1$. Using (3),

$$\begin{aligned} &h^{def}(0)e(-1) - e(-1)h^{def}(0) + h(0)e^{def}(-1) - e^{def}(-1)h(0) \\ &= (h^{def}(0)e(-1)\mathbf{1})_{0-1-0} + (h(0)e(-1)\mathbf{1})_{0-1-0}^{def} = 2e^{def}(-1). \end{aligned} \quad (6)$$

Apply these modes on $e(-1)^{i-1}\mathbf{1}$, we see that

$$\begin{aligned} &h^{def}(0)e(-1)^i\mathbf{1} - e(-1)h^{def}(0)e(-1)^{i-1}\mathbf{1} + h(0)e^{def}(-1)e(-1)^{i-1}\mathbf{1} - e^{def}(-1)h(0)e(-1)^{i-1}\mathbf{1} \\ &= 2e^{def}(-1)e(-1)^{i-1}\mathbf{1}. \end{aligned} \quad (7)$$

Since there are no relations in conformal weights at most k , the coboundary construction Section 3 of [KQ] applies verbatim in this range. Replacing Y_1 by an equivalent cocycle, we may therefore assume that the Y_1 -modes in this range agree with the particular solution constructed in Section 4.1 of [KQ]. In particular, for every $j \leq k-1$, the mode formula in Section 4.1 (4) of [KQ] states that $e(-1)^{def}e(-1)^j\mathbf{1}$ may be computed by

$$\frac{1}{2} \sum_{\beta=1}^j e(-1)^{\beta-1} \sum_{\alpha \geq 0} \binom{-1}{\alpha} \left((e(\alpha)e(-1)\mathbf{1})_{-1+1-\alpha}^{def} + (e^{def}(\alpha)e(-1)\mathbf{1})_{-1+1-\alpha} \right) e(-1)^{j-\beta}\mathbf{1}$$

Since both $e(\alpha)e(-1)\mathbf{1}$ and $e^{def}(\alpha)e(-1)\mathbf{1}$ vanishes for every $\alpha \geq 0$, we conclude that for every $j \leq k-1$,

$$e(-1)^{def}e(-1)^j\mathbf{1} = 0. \quad (8)$$

With (8) and the induction hypothesis, (7) simplifies to

$$h^{def}(0)e(-1)^i\mathbf{1} = 0$$

This proves the claim.

4.3. The terms with $e^{def}(-1)$. We first show that (8) also holds for $i = k$. Apply (6) to $e(-1)^k \mathbf{1}$.

$$\begin{aligned} & h^{def}(0)e(-1)^{k+1} \mathbf{1} - e(-1)h^{def}(0)e(-1)^k \mathbf{1} + h(0)e^{def}(-1)e(-1)^k \mathbf{1} - e^{def}(-1)h(0)e(-1)^k \mathbf{1} \\ &= 2e^{def}(-1)e(-1)^k \mathbf{1}. \end{aligned}$$

We analyze the left-hand-side. The first term vanishes because $e(-1)^{k+1} \mathbf{1} = 0$. The second term vanishes by the conclusion in Section 4.2. The fourth term is $2ke^{def}(-1)e(-1)^k \mathbf{1}$. Moving the fourth term to the right-hand-side results in

$$h(0)e^{def}(-1)e(-1)^k \mathbf{1} = (2k + 2)e^{def}(-1)e(-1)^k \mathbf{1}$$

So $e^{def}(-1)e(-1)^k \mathbf{1}$ is of conformal weight $k + 1$ and of h -weight $2k + 2$. Clearly, no such nonzero vector can exist. The conclusion then follows. In particular, the term $f(1)e^{def}(-1)e(-1)^{i-1} \mathbf{1}$ in (5) vanishes.

Another term in (5) with $e^{def}(-1)$ is $e^{def}(-1)f(1)e(-1)^{i-1} \mathbf{1}$. We calculate that

$$\begin{aligned} f(1)e(-1)^{i-1} \mathbf{1} &= \sum_{j=0}^{i-2} e(-1)^j [f(1), e(-1)] e(-1)^{i-2-j} \mathbf{1} \\ &= \sum_{j=0}^{i-2} e(-1)^j (-h(0) + k) e(-1)^{i-2-j} \mathbf{1} \\ &= \sum_{j=0}^{i-2} e(-1)^j (-2(i-2-j) + k) e(-1)^{i-2-j} \mathbf{1} \\ &= ((-2i + 4 + k)(i-1) + (i-2)(i-1)) e(-1)^{i-2} \mathbf{1} \\ &= (i-1)(k-i+2) e(-1)^{i-2} \mathbf{1}. \end{aligned} \tag{9}$$

Therefore from (8), we conclude that $e^{def}(-1)f(1)e(-1)^{i-1} \mathbf{1} = 0, 1 \leq i \leq k + 1$.

4.4. Proof of $c = 0$. We are now ready to finish the proof. Since $e(-1)^{k+1} \mathbf{1} = 0$, with the results in Section 4.2 and 4.3, the recursion (5) with $i = k + 1$ yields

$$0 = e(-1)f^{def}(1)e(-1)^k \mathbf{1} + ce(-1)^k \mathbf{1}$$

From the recursion (5) with $i = k$, we may obtain

$$0 = e(-1)^2 f^{def}(1)e(-1)^{k-1} \mathbf{1} + 2ce(-1)^k \mathbf{1}$$

Repeatedly apply (5) with $i = k - 1, \dots, 1$, we obtain

$$0 = e(-1)^{k+1} f^{def}(1) \mathbf{1} + (k+1) \cdot ce(-1)^k \mathbf{1}$$

Since $f^{def}(1) \mathbf{1} = 0$, we obtain that

$$(k+1) \cdot ce(-1)^k \mathbf{1} = 0$$

Therefore, $c = 0$. We then conclude the following theorem.

Theorem 4.1. A simple affine vertex operator algebra with positive integer level admits only trivial deformations.

5. AFFINE $\mathfrak{sl}_2$ WITH ADMISSIBLE LEVEL $-4/3$

We now study the case $V = L_{-4/3}(\mathfrak{sl}_2)$. It is known (see [A] for details) that there is a singular vector of conformal weight 3 that is of the form

$$s = -48h(-1)e(-2)\mathbf{1} + 36e(-1)^2f(-1)\mathbf{1} - 6h(-2)e(-1)\mathbf{1} + 9h(-1)^2e(-1)\mathbf{1} + 80e(-3)\mathbf{1}$$

The relevant basis elements in $V^{-4/3}(\mathfrak{g})$ are

$$h(-1)e(-2)\mathbf{1}, e(-1)^2f(-1)\mathbf{1}, h(-2)e(-1)\mathbf{1}, h(-1)^2e(-1)\mathbf{1}, e(-3)\mathbf{1}.$$

The strategy is basically the same as that in Section 4, except that we need to consider both the $f^{def}(1)s$ and $h^{def}(1)s$ to obtain enough constraints for proving $c = 0$.

5.1. $f^{def}(1)$ -action on the singular vector. We note that the formulas in Section 4.1 (4) of [KQ] stay valid for every element of conformal weight 1 and 2. Using those formulas, together with the commutator condition as in Section 4.1, we compute that

$$\begin{aligned} f^{def}(1)h(-1)e(-2)\mathbf{1} &= -f(1)h^{def}(-1)e(-2)\mathbf{1}. \\ f^{def}(1)e(-1)^2f(-1)\mathbf{1} &= 2c \cdot e(-1)f(-1)\mathbf{1} - f(1)e^{def}(-1)e(-1)f(-1)\mathbf{1}. \\ f^{def}(1)h(-2)e(-1)\mathbf{1} &= c \cdot h(-2)\mathbf{1} - f(1)h^{def}(-2)e(-1)\mathbf{1}. \\ f^{def}(1)h(-1)^2e(-1)\mathbf{1} &= c \cdot h(-1)^2\mathbf{1} - f(1)h^{def}(-1)h(-1)e(-1)\mathbf{1}. \\ f^{def}(1)e(-3)\mathbf{1} &= 0. \end{aligned}$$

Therefore,

$$\begin{aligned} f^{def}(1)s &= -48f^{def}(1)h(-1)e(-2)\mathbf{1} + 36f^{def}(1)e(-1)^2f(-1)\mathbf{1} - 6f^{def}(1)h(-2)e(-1)\mathbf{1} \\ &\quad + 9f^{def}(1)h(-1)^2e(-1)\mathbf{1} + 80f^{def}(1)e(-3)\mathbf{1} \\ &= 48f(1)h^{def}(-1)e(-2)\mathbf{1} + 72c \cdot e(-1)f(-1)\mathbf{1} - 36f(1)e^{def}(-1)e(-1)f(-1)\mathbf{1} - 6c \cdot h(-2)\mathbf{1} \\ &\quad + 6 \cdot f(1)h^{def}(-2)e(-1)\mathbf{1} + 9 \cdot c \cdot h(-1)^2\mathbf{1} - 9 \cdot f(1)h^{def}(-1)h(-1)e(-1)\mathbf{1} \end{aligned}$$

5.2. $h^{def}(1)$ -actions on the singular vector. With a similar process as in Section 5.1, we compute that

$$\begin{aligned} h^{def}(1)h(-1)e(-2)\mathbf{1} &= -h(1)h^{def}(-1)e(-2)\mathbf{1} + 2c \cdot e(-2)\mathbf{1}. \\ h^{def}(1)e(-1)e(-1)f(-1)\mathbf{1} &= -h(1)e^{def}(-1)e(-1)f(-1)\mathbf{1} \\ h^{def}(1)h(-2)e(-1)\mathbf{1} &= -h(1)h^{def}(-2)e(-1)\mathbf{1}. \\ h^{def}(1)h(-1)h(-1)e(-1)\mathbf{1} &= 4c \cdot h(-1)e(-1)\mathbf{1} - h(1)h^{def}(-1)h(-1)e(-1)\mathbf{1}. \\ h^{def}(1)e(-3)\mathbf{1} &= 0. \end{aligned}$$

Therefore,

$$\begin{aligned}
h^{def}(1)s &= -48h^{def}(1)h(-1)e(-2)\mathbf{1} + 36h^{def}(1)e(-1)^2f(-1)\mathbf{1} - 6h^{def}(1)h(-2)e(-1)\mathbf{1} \\
&\quad + 9h^{def}(1)h(-1)^2e(-1)\mathbf{1} + 80h^{def}(1)e(-3)\mathbf{1} \\
&= 48h(1)h^{def}(-1)e(-2)\mathbf{1} - 96c \cdot e(-2)\mathbf{1} - 36h(1)e^{def}(-1)e(-1)f(-1)\mathbf{1} \\
&\quad + 6h(1)h^{def}(-2)e(-1)\mathbf{1} + 36c \cdot h(-1)e(-1)\mathbf{1} - 9h(1)h^{def}(-1)h(-1)e(-1)\mathbf{1}.
\end{aligned}$$

5.3. Proof of $c = 0$. Note that since $h^{def}(-1)e(-1)\mathbf{1} = 0$, by D -commutator formula, it is necessary that

$$h^{def}(-2)e(-1)\mathbf{1} + h^{def}(-1)e(-2)\mathbf{1} = 0.$$

Therefore, it suffices to set $h^{def}(-1)e(-2)\mathbf{1}, e^{def}(-1)e(-1)f(-1)\mathbf{1}, h^{def}(-1)h(-1)e(-1)\mathbf{1}$ as generic elements of the conformal weight 3. Also, to determine c it suffices to consider only those terms with Cartan weight 2. So we set

$$\begin{aligned}
h^{def}(-1)e(-2)\mathbf{1} &= a_1e(-3)\mathbf{1} + a_2e(-2)h(-1)\mathbf{1} + a_3e(-1)h(-2)\mathbf{1} \\
&\quad + a_4e(-1)h(-1)^2\mathbf{1} + a_5e(-1)^2f(-1)\mathbf{1}, \\
h^{def}(-1)h(-1)e(-1)\mathbf{1} &= b_1e(-3)\mathbf{1} + b_2e(-2)h(-1)\mathbf{1} + b_3e(-1)h(-2)\mathbf{1} \\
&\quad + b_4e(-1)h(-1)^2\mathbf{1} + b_5e(-1)^2f(-1)\mathbf{1}. \\
e^{def}(-1)e(-1)f(-1)\mathbf{1} &= c_1e(-3)\mathbf{1} + c_2e(-2)h(-1)\mathbf{1} + c_3e(-1)h(-2)\mathbf{1} \\
&\quad + c_4e(-1)h(-1)^2\mathbf{1} + c_5e(-1)^2f(-1)\mathbf{1},
\end{aligned}$$

Plug them into the formula of $f^{def}(1)s$ obtained in Section 5.1 and $h^{def}(1)s$ obtained in Section 5.2, and set them to be zero, we obtain the following five equations concerning the structural constants:

$$0 = 84a_3 + 168a_4 - 28a_5 - 18b_3 - 36b_4 + 6b_5 - 72c_3 - 144c_4 + 24c_5 + 12c \quad (10)$$

$$0 = -42a_2 - 56a_4 + 9b_2 + 12b_4 + 36c_2 + 48c_4 + 9c \quad (11)$$

$$0 = -42a_1 - 56a_3 + 9b_1 + 12b_3 + 36c_1 + 48c_3 - 6c \quad (12)$$

$$0 = 84a_2 - 560a_4 + 168a_5 - 18b_2 + 120b_4 - 36b_5 - 72c_2 + 480c_4 - 144c_5 + 36c \quad (13)$$

$$\begin{aligned}
0 &= 84a_1 - 280a_2 - 168a_3 + 784a_4 - 336a_5 - 18b_1 + 60b_2 + 36b_3 - 168b_4 + 72b_5 \\
&\quad - 72c_1 + 240c_2 + 144c_3 - 672c_4 + 288c_5 - 96c
\end{aligned} \quad (14)$$

We apply Gaussian elimination. The operation

$$(14) + 2(12) - \frac{70 \cdot 4}{42}(11) + \frac{280}{84}(10)$$

eliminates a_1, a_2, a_3 from (14) and results in

$$\frac{4}{3}(1288a_4 - 322a_5 - 276b_4 + 69b_5 - 96c - 1104c_4 + 276c_5) = 0 \quad (15)$$

To eliminate a_4 in (15), we first eliminate a_2 from (13) by the operation

$$(13) + 2(11),$$

which results in

$$-6(112a_4 - 28a_5 - 24b_4 + 6b_5 - 9c - 96c_4 + 24c_5) = 0 \quad (16)$$

Then, the operation

$$(15) + \frac{1288 \cdot 4}{112 \cdot 6 \cdot 3}(16)$$

results in

$$10c = 0 \Rightarrow c = 0.$$

Theorem 5.1. The simple affine VOA $L_{-4/3}(sl_2)$ admits only trivial deformations.

AUTHORS' STATEMENT

The manuscript has no associated data. There exists no conflict of interest. We used ChatGPT 5.6-Pro to detect the gaps and mistakes in the previous version. The proof of Theorem 3.1 in the current version is contributed by ChatGPT 5.6-Pro.

REFERENCES

- [A] D. Adamović, A construction of admissible $A_1(1)$ -modules of level $-4/3$, *J. Pure and Applied Alg.* **196** (2005) 119–134
- [FZ] Igor Frenkel and Yongchang Zhu, Vertex operator algebras associated to representations of affine and Virasoro algebras, *Duke Math. J.* **66** (1992), no. 1, 123–168.
- [H1] Yi-Zhi Huang, A cohomology theory of grading-restricted vertex algebras, *Comm. Math. Phys.* **327** (2014), 279–307.
- [H2] Yi-Zhi Huang, First and second cohomologies of grading-restricted vertex algebras, *Comm. Math. Phys.* **327** (2014), 261–278.
- [H3] Yi-Zhi Huang, Addendum to the proof of Proposition 1.4 in [H2], <https://sites.math.rutgers.edu/~yzhuang/rci/math/papers/2nd-coh-addendum.pdf>
- [H4] Yi-Zhi Huang, Some open problems in mathematical two-dimensional conformal field theory, in: *Proceedings of the Conference on Lie Algebras, Vertex Operator Algebras, and Related Topics, held at University of Notre Dame*, Notre Dame, Indiana, August 14–18, 2015, ed. K. Barron, E. Jurisich, H. Li, A. Milas, K. C. Misra, Contemp. Math, Vol. 695, American Mathematical Society, Providence, RI, 2017, 123–138.
- [GK] Maria Gorelik and Victor Kac, On simplicity of vacuum modules. *Advances in Mathematics*, **211**(2): 621–677, 2007.
- [KQ] Vladimir Kovalchuk and Fei Qi, First-order deformations of freely generated vertex algebras, arXiv:2408.16309.
- [KW] Victor Kac and Minoru Wakimoto, Modular invariant representations of infinite dimensional Lie algebras and superalgebras, *Proc. Natl. Acad. Sci. USA* **85** (1988), 4956–4960.
- [LL] James Lepowsky and Haisheng Li, *Introduction to Vertex Operator Algebras and Their Representations*, Progress in Math., Vol. 227, Birkhäuser, Boston, 2003.
- [L] Bong H. Lian, On the classification of simple vertex operator algebras, *Comm. Math. Phys.* **163** (1994), no. 2, 307–357.

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF DENVER, DENVER, CO 80210, USA

E-mail address: andrew.linshaw@du.edu

SCHOOL OF MATHEMATICS (ZHUHAI), SUN YAT-SEN UNIVERSITY, ZHUHAI, GUANGDONG, CHINA

E-mail address: qifei@mail.sysu.edu.cn | fei.qi.math.phys@gmail.com